

**POSTECH 이성익 교수의
양자 세계에 관한 강연
- 2장 -**

편집 도우미: POSTECH 학부생 정운영

Chapter 2

Wave Particle Duality, Probability, and the Schrödinger Equation

$$\psi(x, t) = \int_{-\infty}^{\infty} A(k) e^{i(kx - \omega t)} dk$$

t=0에 서,

$$\psi(x) = \int_{-\infty}^{\infty} g(k) e^{ikx} dk$$

이 때 $g(k) = e^{-\alpha(k-k_0)^2}$ 로 주어지면

$$\psi(x) = \int_{-\infty}^{\infty} e^{-\alpha(k-k_0)^2} \cdot e^{ikx} dk = \int_{-\infty}^{\infty} e^{-\alpha\left(k^2 - 2k_0k + k_0^2 - \frac{ikx}{\alpha}\right)} dk$$

$$\begin{aligned} -\alpha\left(k^2 - 2k_0k + k_0^2 - \frac{ikx}{\alpha}\right) &= -\alpha\left(k^2 - 2\left(k_0 + \frac{ix}{2\alpha}\right)k + k_0^2\right) \\ &= -\alpha\left(k^2 - 2\left(k_0 + \frac{ix}{2\alpha}\right)k + k_0^2\right) \\ &= -\alpha\left(k^2 - 2\left(k_0 + \frac{ix}{2\alpha}\right)k + \left(k_0 + \frac{ix}{2\alpha}\right)^2 + k_0^2 - k_0^2 - \frac{ik_0x}{\alpha} + \frac{x^2}{4\alpha^2}\right) \\ &= -\alpha\left(\left(k - k_0 - \frac{ix}{2\alpha}\right)^2 + \frac{x^2}{4\alpha^2} - \frac{ik_0x}{\alpha}\right) \end{aligned}$$

$$\psi(x) = \sqrt{\frac{\pi}{\alpha}} \cdot e^{-\alpha\left(\frac{x^2}{4\alpha^2} - \frac{ik_0x}{\alpha}\right)} = \sqrt{\frac{\pi}{\alpha}} \cdot e^{ik_0x} \cdot e^{-\frac{x^2}{4\alpha}}$$

Free space에 있는 wave는 시간이 지남에 따라 퍼져나갈 것이다.

$$\psi(x, t) = \int_{-\infty}^{\infty} A(k) e^{ik(x-vt)} dk = f(x-vt)$$

$$\left(e^{ikx-ivt} = e^{i\frac{2\pi}{\lambda}x - i2\pi vt} = e^{2\pi i\left(\frac{x}{\lambda} - vt\right)} = e^{2\pi i\left(\frac{x}{\lambda} - \frac{v}{\lambda}t\right)} = e^{ik(x-vt)} \right)$$

$$\psi(x, t) = \int_{-\infty}^{\infty} A(k) e^{i(kx-w(k)t)} dk$$

$$A(k) = e^{-\alpha(k-k_0)^2} \text{로 주어지면,}$$

$$\psi(x, t) = \int_{-\infty}^{\infty} e^{-\alpha(k-k_0)^2} e^{i(kx-w(k)t)} dk$$

$$\begin{aligned} w(k) &\cong w(k_0) + (k-k_0) \left(\frac{dw}{dk} \right)_{k=k_0} + \frac{1}{2} (k-k_0)^2 \left(\frac{d^2w}{dk^2} \right)_{k=k_0} \\ &= w(k_0) + v_g (k-k_0) + \beta (k-k_0)^2 \quad k-k_0 = k' \end{aligned}$$

$$\begin{aligned} \psi(x, t) &= \int_{-\infty}^{\infty} e^{-\alpha(k-k_0)^2} e^{ikx} e^{-iw(k)t} dk = \int_{-\infty}^{\infty} e^{-\alpha k'^2} e^{ik'_x} e^{ik_0x} e^{-iw(k_0)t} e^{-iv_g k't} e^{-i\beta k'^2 t} dk' \\ &= \int_{-\infty}^{\infty} e^{-\alpha k'^2} e^{ik'_x} e^{-i(v_g k' + \beta k'^2)t} e^{ik_0x} e^{-iw(k_0)t} dk' \\ &= e^{ik_0x} e^{-iw(k_0)t} \int_{-\infty}^{\infty} e^{-\alpha k'^2} e^{-i\beta t k'^2} e^{ik'(x-v_g t)} dk' \\ &= e^{i(k_0x-w(k_0)t)} \int_{-\infty}^{\infty} e^{ik'(x-v_g t)} e^{-(\alpha+i\beta t)k'^2} dk' \\ &= e^{i(k_0x-w(k_0)t)} \int_{-\infty}^{\infty} e^{-\alpha+i\beta t \left(k'^2 - \frac{ik'(x-v_g t)}{\alpha+i\beta t} \right)} dk' \\ &= e^{i(k_0x-w(k_0)t)} \int_{-\infty}^{\infty} e^{-\alpha+i\beta t \left(k' - \frac{i(x-v_g t)}{2(\alpha+i\beta t)} \right)^2 - \frac{(x-v_g t)^2}{4(\alpha+i\beta t)}} dk' \end{aligned}$$

$$\begin{aligned}
&= e^{i(k_0x-w(k_0)t)} e^{-\frac{(x-v_g t)^2}{4(\alpha+i\beta t)}} \int_{-\infty}^{\infty} e^{-i(\alpha+i\beta t) \left(k - \frac{i(x-v_g t)}{2(\alpha+i\beta t)} \right)^2} dk \\
&= e^{i(k_0x-w(k_0)t)} \left(\frac{\pi}{\alpha+i\beta t} \right)^{\frac{1}{2}} e^{-\frac{(x-v_g t)^2}{4(\alpha+i\beta t)}} \\
\psi^*(x,t) &= e^{-i(k_0x-w(k_0)t)} \left(\frac{\pi}{\alpha-i\beta t} \right)^{\frac{1}{2}} e^{-\frac{(x-v_g t)^2}{4(\alpha-i\beta t)}} \\
|\psi(x,t)|^2 &= \left(\frac{\pi^2}{\alpha^2 + \beta^2 t^2} \right)^{\frac{1}{2}} \cdot e^{-\frac{\alpha(x-v_g t)^2}{2(\alpha^2 + \beta^2 t^2)}}
\end{aligned}$$

Schrödinger Equation

free space에 대해,

$$p = \frac{h}{\lambda} = h \frac{k}{2\pi} = \hbar k \qquad E = h\nu = \frac{h}{2\pi} 2\pi\nu = \hbar\omega$$

$$\psi(x,t) = \int_{-\infty}^{\infty} A(k) e^{i(kx-wt)} dk = \int_{-\infty}^{\infty} \phi(p) e^{i(px-Et)/\hbar} dp$$

$$\begin{aligned}
i\hbar \frac{\partial \psi(x,t)}{\partial t} &= \int_{-\infty}^{\infty} \phi(p) \left(-\frac{iE}{\hbar} \right) (i\hbar) e^{i(px-Et)/\hbar} dp = E \int_{-\infty}^{\infty} \phi(p) e^{i(px-Et)/\hbar} dp \\
&= E\psi(x,t) = \frac{p^2}{2m} \psi(x,t) = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x,t)}{\partial x^2}
\end{aligned}$$

$$\therefore i\hbar \frac{\partial \psi(x,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x,t)}{\partial x^2}$$

→ given by nature

Potential이 있으면,

$$E = \frac{p^2}{2m} + V$$
$$i\hbar \frac{\partial \psi(x,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x,t)}{\partial x^2} + V\psi(x,t)$$

Max Born,

$$|\psi|^2 = \psi^* \cdot \psi \rightarrow \text{probability}$$

$$P(x,t)dx = |\psi(x,t)|^2 dx$$

$$\int_{-\infty}^{\infty} P(x,t)dx = 1$$

Uncertainty Principle

$$\Delta x \cdot \Delta p \geq \frac{\hbar}{2}$$

이 것을 Bohr의 수소 원자모형에 응용하면,

$$r \cdot p \sim \hbar$$

$$E = \frac{p^2}{2m} - \frac{e^2}{4\pi\epsilon_0 r} \simeq \frac{\hbar^2}{2mr^2} - \frac{e^2}{4\pi\epsilon_0 r}$$

$$\frac{\partial E}{\partial r} = -\frac{\hbar^2}{mr^3} + \frac{e^2}{4\pi\epsilon_0 r^2} = 0$$

$$\frac{\hbar^2}{mr} = \frac{e^2}{4\pi\epsilon_0}$$

$$r = \frac{4\pi\epsilon_0 \hbar^2}{me^2} = \frac{\hbar}{mc} \frac{4\pi\epsilon_0 \hbar c}{e^2} = \frac{\hbar}{mc} \frac{1}{\alpha}$$

$$E \simeq \frac{\hbar^2}{2mr^2} - \frac{e^2}{4\pi\epsilon_0 r} = \frac{\hbar^2}{2m} \frac{m^2 c^2 \alpha^2}{\hbar^2} - \frac{e^2}{4\pi\epsilon_0} \frac{m c \alpha}{\hbar}$$

$$= \frac{m c^2}{2} \alpha^2 - m c^2 \alpha^2 = -\frac{m c^2}{2} \alpha^2$$

Conjugate Equation

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi \qquad \frac{\partial \psi}{\partial t} = \frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{i\hbar} V\psi$$

$$-i\hbar \frac{\partial \psi^*}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi^*}{\partial x^2} + V\psi^* \qquad \frac{\partial \psi^*}{\partial t} = -\frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} - \frac{1}{i\hbar} V\psi^*$$

$$\frac{\partial}{\partial t} P(x,t) = \frac{\partial}{\partial t} (\psi^*(x,t) \cdot \psi(x,t))$$

$$= \frac{\partial \psi^*(x,t)}{\partial t} \psi(x,t) + \frac{\partial \psi(x,t)}{\partial t} \psi^*(x,t)$$

$$= -\frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} \psi - \frac{1}{i\hbar} V\psi^* \psi + \frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} \psi^* + \frac{1}{i\hbar} V\psi \psi^*$$

$$= -\frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} \psi + \frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} \psi^* = \frac{i\hbar}{2m} \left(\psi^* \frac{\partial^2 \psi}{\partial x^2} - \psi \frac{\partial^2 \psi^*}{\partial x^2} \right)$$

$$= \frac{i\hbar}{2m} \frac{\partial}{\partial x} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right)$$

$$\frac{\partial P}{\partial t} + \nabla \cdot \vec{j} = 0 \quad \rightarrow \text{continuity equation}$$

$$j = \frac{-i\hbar}{2m} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right)$$

$$\langle f(x) \rangle = \int_{-\infty}^{\infty} f(x) P(x,t) dx = \int_{-\infty}^{\infty} \psi^*(x,t) f(x) \psi(x,t) dx$$

$$\langle p \rangle = m \frac{d\langle x \rangle}{dt}$$

$$\begin{aligned} &= m \frac{d}{dt} \int_{-\infty}^{\infty} \psi^*(x,t) x \psi(x,t) dx = m \int_{-\infty}^{\infty} \left(\frac{d\psi^*}{dt} x \psi + \frac{d\psi}{dt} x \psi^* \right) dx \\ &= m \int_{-\infty}^{\infty} \left(\left(-\frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} - \frac{1}{i\hbar} V \psi^* \right) x \psi + \left(\frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{i\hbar} V \psi \right) x \psi^* \right) dx \\ &= m \int_{-\infty}^{\infty} \left(-\frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} x \psi - \frac{1}{i\hbar} V \psi^* x \psi + \frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} x \psi^* + \frac{1}{i\hbar} V \psi x \psi^* \right) dx \\ &= m \int_{-\infty}^{\infty} \left(-\frac{i\hbar}{2m} \frac{\partial^2 \psi^*}{\partial x^2} x \psi + \frac{i\hbar}{2m} \frac{\partial^2 \psi}{\partial x^2} x \psi^* \right) dx \\ &= m \int_{-\infty}^{\infty} \frac{\hbar}{2mi} \left(\frac{\partial^2 \psi^*}{\partial x^2} x \psi - \frac{\partial^2 \psi}{\partial x^2} x \psi^* \right) dx \\ &= \frac{\hbar}{2i} \int_{-\infty}^{\infty} \left(\frac{\partial^2 \psi^*}{\partial x^2} x \psi - \frac{\partial^2 \psi}{\partial x^2} x \psi^* \right) dx \end{aligned}$$

$$\begin{aligned} &\frac{\partial^2 \psi^*}{\partial x^2} x \psi - \frac{\partial^2 \psi}{\partial x^2} x \psi^* \\ &= \frac{\partial}{\partial x} \left(\frac{\partial \psi^*}{\partial x} x \psi - \frac{\partial \psi}{\partial x} x \psi^* \right) - \frac{\partial \psi^*}{\partial x} \left(\psi + x \frac{\partial \psi}{\partial x} \right) + \frac{\partial \psi}{\partial x} \left(\psi^* + x \frac{\partial \psi^*}{\partial x} \right) \\ &= \frac{\partial}{\partial x} \left(\frac{\partial \psi^*}{\partial x} x \psi - \frac{\partial \psi}{\partial x} x \psi^* \right) - \frac{\partial \psi^*}{\partial x} \psi + \frac{\partial \psi}{\partial x} \psi^* \\ &= \frac{\hbar}{2i} \int_{-\infty}^{\infty} \left(\frac{\partial}{\partial x} \left(\frac{\partial \psi^*}{\partial x} x \psi - \frac{\partial \psi}{\partial x} x \psi^* \right) + \frac{\partial \psi}{\partial x} \psi^* - \psi \frac{\partial \psi^*}{\partial x} \right) dx \\ &= \frac{\hbar}{2i} \int_{-\infty}^{\infty} \left(-\frac{\partial \psi^*}{\partial x} \psi + \psi^* \frac{\partial \psi}{\partial x} \right) dx = \frac{\hbar}{2i} \int_{-\infty}^{\infty} \left(\psi^* \frac{\partial \psi}{\partial x} + \psi \frac{\partial \psi^*}{\partial x} \right) dx \end{aligned}$$

$$= \frac{\hbar}{i} \int_{-\infty}^{\infty} \psi^* \frac{\partial \psi}{\partial x} dx = \int_{-\infty}^{\infty} \psi^* \frac{\hbar}{i} \frac{\partial}{\partial x} \psi dx$$

$$\therefore \langle p \rangle = \int_{-\infty}^{\infty} \psi^*(x,t) \cdot \frac{\hbar}{i} \frac{\partial}{\partial x} \cdot \psi(x,t) dx \quad p = \frac{\hbar}{i} \frac{\partial}{\partial x}$$

$$\langle f(p) \rangle = \int_{-\infty}^{\infty} \psi^*(x,t) \cdot f\left(\frac{\hbar}{i} \frac{\partial}{\partial x}\right) \cdot \psi(x,t) dx$$

t=0 일 때 ,

$$\begin{aligned} \psi(x) &= \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \phi(p) e^{ipx/\hbar} dp \\ &= \sqrt{\frac{\hbar}{2\pi}} \int_{-\infty}^{\infty} \phi(\hbar k) e^{ikx} dk \end{aligned}$$

$$\phi(\hbar k) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \psi(x) e^{-ikx} dx$$

$$\phi(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \psi(x) e^{-ipx/\hbar} dx$$

$$\begin{aligned} \int_{-\infty}^{\infty} \phi^*(p) \cdot \phi(p) dp &= \int_{-\infty}^{\infty} \phi^*(p) \left(\frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \psi(x) e^{-ipx/\hbar} dx \right) dp \\ &= \int_{-\infty}^{\infty} \psi(x) \left(\frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \phi^*(p) e^{-ipx/\hbar} dp \right) dx \\ &= \int_{-\infty}^{\infty} \psi(x) \cdot \psi^*(x) dx = 1 \end{aligned}$$

$$\therefore \int_{-\infty}^{\infty} \phi^*(p) \cdot \phi(p) dp = 1$$

$$\begin{aligned} \langle p \rangle &= \int_{-\infty}^{\infty} \psi^*(x) \cdot \frac{\hbar}{i} \frac{\partial}{\partial x} \cdot \psi(x) dx \\ &= \int_{-\infty}^{\infty} \psi^*(x) \cdot \frac{\hbar}{i} \frac{\partial}{\partial x} \cdot \left(\frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \phi(p) e^{ipx/\hbar} dp \right) dx \\ &= \int_{-\infty}^{\infty} \psi^*(x) \cdot \frac{\hbar}{i} \frac{ip}{\hbar} \left(\frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \phi(p) e^{ipx/\hbar} dp \right) dx \\ &= \int_{-\infty}^{\infty} \phi(p) \cdot p \cdot \left(\frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \psi^*(x) e^{ipx/\hbar} dx \right) dp \\ &= \int_{-\infty}^{\infty} \phi(p) \cdot p \cdot \phi^*(p) dp \end{aligned}$$

$$\therefore \langle p \rangle = \int_{-\infty}^{\infty} \phi^*(p) \cdot p \cdot \phi(p) dp$$

$$\langle f(x) \rangle = \int_{-\infty}^{\infty} \phi^*(p) \cdot f \left(i\hbar \frac{\partial}{\partial p} \right) \cdot \phi(p) dp \quad x = i\hbar \frac{\partial}{\partial p}$$

Generally, operators do not commute

$$[A, B] = AB - BA$$

$$\begin{aligned} [p, x]\psi(x) &= \left[\frac{\hbar}{i} \frac{\partial}{\partial x}, x \right] \psi(x) \\ &= \left(\frac{\hbar}{i} \frac{\partial}{\partial x} x - x \frac{\hbar}{i} \frac{\partial}{\partial x} \right) \psi(x) \\ &= \frac{\hbar}{i} \frac{\partial}{\partial x} x \psi(x) - x \frac{\hbar}{i} \frac{\partial \psi(x)}{\partial x} \\ &= \frac{\hbar}{i} \psi(x) + x \frac{\hbar}{i} \frac{\partial \psi(x)}{\partial x} - x \frac{\hbar}{i} \frac{\partial \psi(x)}{\partial x} \\ &= -i\hbar \psi(x) \end{aligned}$$

$$\therefore [p, x] = -i\hbar$$

$$\begin{aligned} \langle p \rangle - \langle p \rangle^* &= \int_{-\infty}^{\infty} \psi^*(x) \cdot \frac{\hbar}{i} \frac{\partial}{\partial x} \cdot \psi(x) dx - \int_{-\infty}^{\infty} \psi(x) \cdot \left(-\frac{\hbar}{i} \frac{\partial}{\partial x} \right) \cdot \psi^*(x) dx \\ &= \frac{\hbar}{i} \int_{-\infty}^{\infty} \psi^*(x) \cdot \frac{\partial}{\partial x} \psi(x) + \psi(x) \cdot \frac{\partial}{\partial x} \psi^*(x) dx \\ &= \frac{\hbar}{i} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} (\psi^*(x) \cdot \psi(x)) dx = \frac{\hbar}{i} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} |\psi(x)|^2 dx \\ &= \frac{\hbar}{i} \left(|\psi(x)|^2 \Big|_{x=\infty} - |\psi(x)|^2 \Big|_{x=-\infty} \right) \\ &= 0 \end{aligned}$$